

Open Problems in the Theory of Macdonald Polynomials

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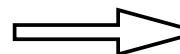
A remarkable determinant

(1,0)	(1,1)
(0,0)	(0,1)



$$\Delta_{2,2}(X, Y) = \det \begin{pmatrix} x_1^0 y_1^0 & x_2^0 y_2^0 & x_3^0 y_3^0 & x_4^0 y_4^0 \\ x_1^0 y_1^1 & x_2^0 y_2^1 & x_3^0 y_3^1 & x_4^0 y_4^1 \\ x_1^1 y_1^0 & x_2^1 y_2^0 & x_3^1 y_3^0 & x_4^1 y_4^0 \\ x_1^1 y_1^1 & x_2^1 y_2^1 & x_3^1 y_3^1 & x_4^1 y_4^1 \end{pmatrix} = \det \begin{pmatrix} 1 & 1 & 1 & 1 \\ y_1 & y_2 & y_3 & y_4 \\ x_1 & x_2 & x_3 & x_4 \\ x_1 y_1 & x_2 y_2 & x_3 y_3 & x_4 y_4 \end{pmatrix}$$

(2,0)		
(1,0)	(1,1)	(1,2)
(0,0)	(0,1)	(0,2)



$$\Delta_{331}(X, Y) = \det \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ y_1 & y_2 & y_3 & y_4 & y_5 & y_5 & y_7 \\ y_1^2 & y_2^2 & y_3^2 & y_4^2 & y_5^2 & y_6^2 & y_7^2 \\ x_1 & x_2 & x_3 & x_4 & x_5 & x_6 & x_7 \\ x_1 y_1 & x_2 y_2 & x_3 y_3 & x_4 y_4 & x_5 y_5 & x_6 y_6 & x_7 y_7 \\ x_1 y_1^2 & x_2 y_2^2 & x_3 y_3^2 & x_4 y_4^2 & x_5 y_5^2 & x_6 y_6^2 & x_7 y_7^2 \\ x_1^2 & x_2^2 & x_3^2 & x_4^2 & x_5^2 & x_6^2 & x_7^2 \end{pmatrix}$$

General definition

If $(p_1, q_1), (p_2, q_2), \dots, (p_n, q_n)$ are the cells of the Ferrers diagram of $\mu \vdash n$ then

$$\Delta_\mu(X, Y) = \det \|x_j^{p_i} y_j^{q_i}\|_{i,j=1}^n$$

next

In case we forgot

Theorem (easy)

For any $\mu \vdash n$ the dimension of the linear span of the derivatives of $\Delta_\mu(X, Y)$ is at most $n!$

In symbols

$$\dim M_\mu[X, Y] \leq n!$$

where

$$M_\mu[X, Y] = L[\delta_{x_1}^{r_1} \delta_{y_1}^{s_1} \delta_{x_2}^{r_2} \delta_{y_2}^{s_2} \cdots \delta_{x_n}^{r_n} \delta_{y_n}^{s_n} \Delta_\mu(X, Y) : r_i, s_i \geq 0]$$

For example [using MAPLE]

```
DDmu([2,1]);
```

$$\begin{bmatrix} 1 & 1 & 1 \\ y1 & y2 & y3 \\ x1 & x2 & x3 \end{bmatrix}$$

```
D21:=det(");
```

$$D21 := y2 x3 - y3 x2 - y1 x3 + y1 x2 + x1 y3 - x1 y2$$

```
diff(D21,x1);  
diff(D21,x3);  
diff(D21,y1);  
diff(D21,y3);  
diff(D21,x3,y2);
```

6 independent derivatives!

how do we get these $n!$ derivatives?

next

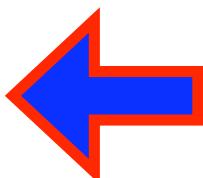
The $n!$ Conjecture

(1990 - 2000)

For $\mu \vdash n$

$$\dim M_\mu[X, Y] = n!$$

Proved by **Mark Haiman** using algebraic geometry

 1,000\$ 
Reward

for an “elementary” proof

“Elementary” means:

By calculus or/and combinatorics

give an algorithm

that produces a “triangular” set of $n!$ derivatives

next

The basic construction

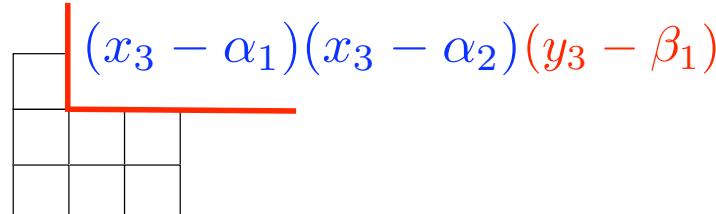
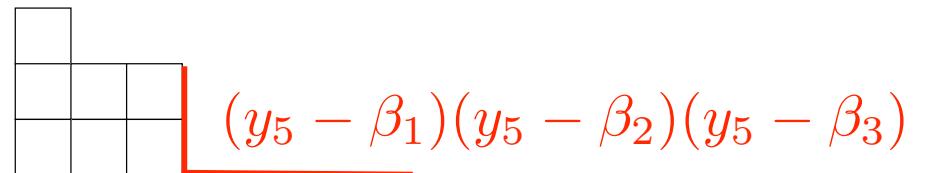
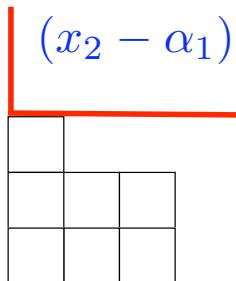
from a tableau to a point in $2n$ dimensional space

$$\begin{array}{c} \alpha_3 \\ \alpha_2 \\ \alpha_1 \\ \beta_1 \beta_2 \beta_3 \end{array} \begin{array}{|c|c|c|} \hline 6 & & \\ \hline 7 & 1 & 4 \\ \hline 2 & 5 & 3 \\ \hline \end{array} \rightarrow \begin{pmatrix} \alpha_2 & \alpha_1 & \alpha_1 & \alpha_2 & \alpha_1 & \alpha_3 & \alpha_2 & \beta_2 & \beta_1 & \beta_3 & \beta_3 & \beta_2 & \beta_1 & \beta_1 \end{pmatrix}_{x_1 \ x_2 \ x_3 \ x_4 \ x_5 \ x_6 \ x_7 \ y_1 \ y_2 \ y_3 \ y_4 \ y_5 \ y_6 \ y_7}$$

some polynomials vanishing at this orbit point

$$x_1 - \alpha_2 \quad x_4 - \alpha_2 \quad x_7 - \alpha_2 \quad y_2 - \beta_1 \quad y_5 - \beta_2$$

Polynomials vanishing at all these tableau points:

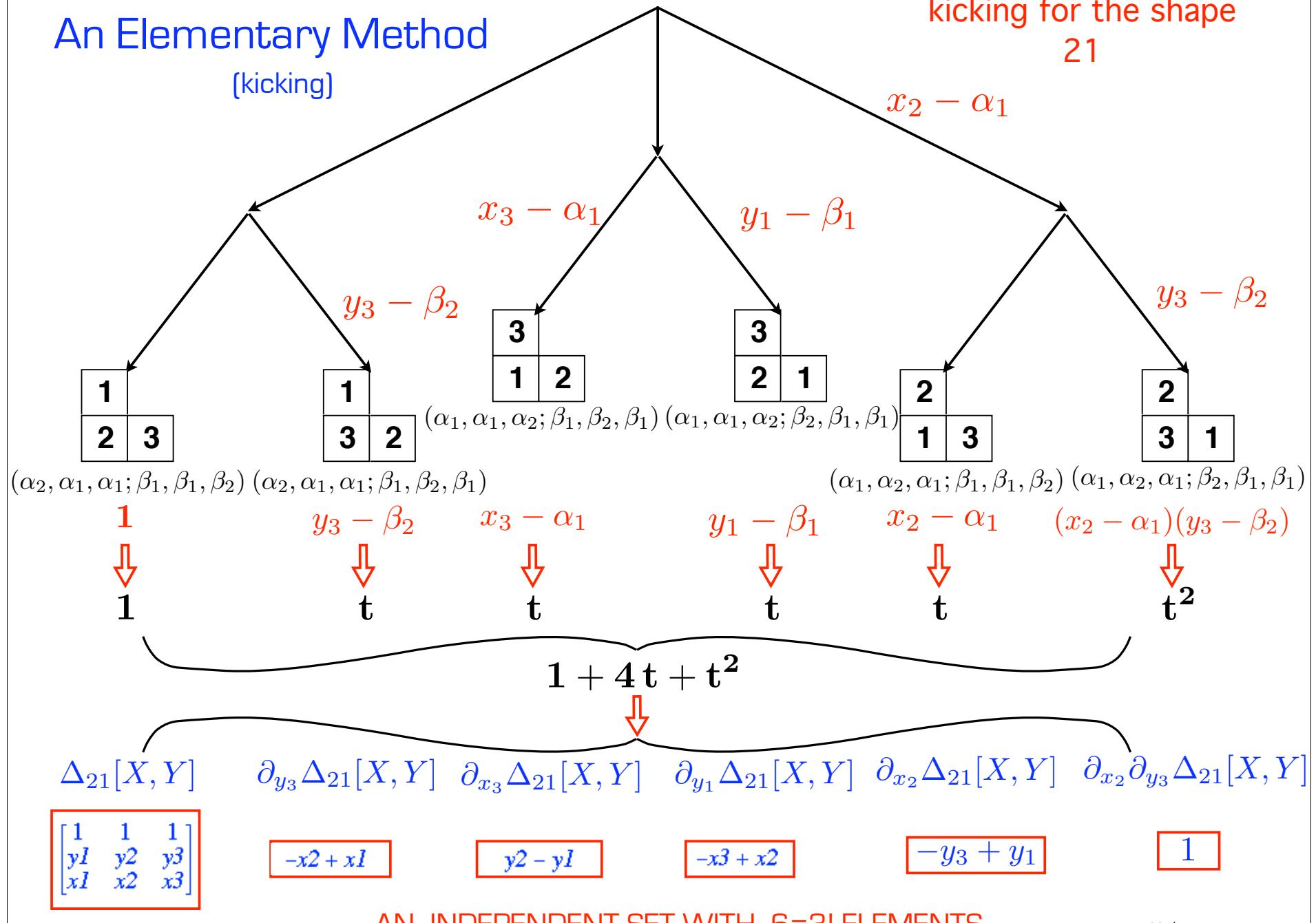


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An Elementary Method

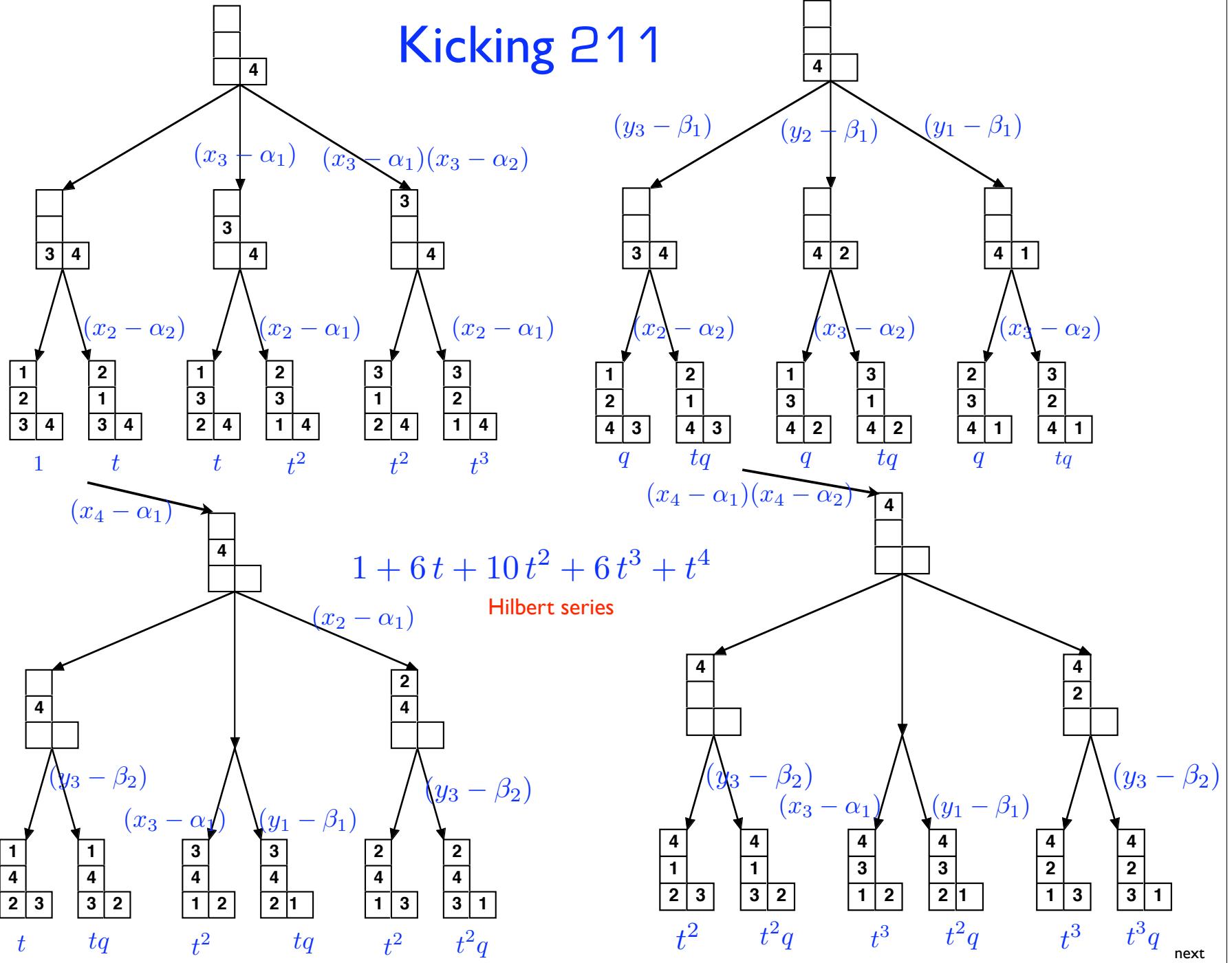
(kicking)

kicking for the shape
21



next

Kicking 211



The Algorithm

Step 1

Fill the Ferrers diagram of μ with $1, 2, \dots, n$ in all possible ways to get the $n!$ tableaux.

Step 2

Construct the corresponding orbit points

5040
In this case

Step 3

Order the orbit points and construct the kicking polynomials.

Step 4

Compute the kicking statistics.

3		
6	1	2
4	7	5

1
12
75
306
807
1319
1319
807
306
75
12
1

Theorem

If the kicking statistic are symmetric then the top components of the kicking polynomials give the $n!$ independent derivatives of $\Delta_\mu(X, Y) = \det \|x_j^{p_i} y_j^{q_i}\|_{i,j=1}^n$

I am planning to advertise kicking
as a replacement

for SUDOKU

next

Recalling Hilbert series

A vector space V is called “*graded*” if and only if

$$V = H_0(V) \oplus H_1(V) \oplus H_2(V) \oplus \cdots \oplus H_m(V) \oplus \cdots$$

If $\dim H_m(V) < \infty$ for all m , we set

$$F_V(t) = \sum_{m \geq 0} t^m \dim H_m(V)$$

the “*Hilbert series*” of V

Our spaces $M_\mu[X, Y]$ are “*bigraded*” that is we have the double decomposition

$$M_\mu[X, Y] = \bigoplus_{r=0}^{n(\mu)} \bigoplus_{s=0}^{n(\mu')} H_{r,s}(M_\mu[X, Y])$$

With $H_{r,s}(M_\mu[X, Y])$ the linear span of derivatives of $\Delta_\mu(x, y)$

that are homogeneous of degree R in x_1, x_2, \dots, x_n and degree S in y_1, y_2, \dots, y_n

Here and after we set

$$F_\mu(q, t) = \sum_{r=0}^{n(\mu)} \sum_{s=0}^{n(\mu')} t^r q^s \dim H_{r,s}(M_\mu[X, Y])$$

Using Maple

`hilb([3,2]);`

$$q^4 t^2 + 4 q^4 t + 4 q^3 t^2 + 5 q^4 + 15 q^3 t + 9 q^2 t^2 + 11 q^3 + 22 q^2 t + 11 q t^2 + 9 q^2 + 15 q t + 5 t^2 + 4 q + 4 t + 1$$

The dimension of

$$H_{2,1}(M_{3,2}[X, Y])$$

2 →	[5	11	9	4	1]
1 →	4	15	22	15	4
0 →	1	4	9	11	5

The dimension of $H_{1,3}(M_{3,2}[X, Y])$

next

Flip Definition

We say that a vector space V of polynomials is a "cone" if it is the linear span of derivatives of a single homogeneous polynomial $\Delta(x)$.

In this case V has an automorphism "Flip" defined by

$$\text{flip } P(x) = P(\partial_x)\Delta(x)$$

Note this is a non-singular linear map of V onto V since if

$$P(x) = Q(\partial_x)\Delta(x) \quad \text{and} \quad P(\partial_x)\Delta(x) = 0$$

then

$$P(\partial_x)P(x) = P(\partial_x)Q(\partial_x)\Delta(x) = Q(\partial_x)P(\partial_x)\Delta(x) = 0$$

Thus

$$\text{flip } P(x) = 0 \implies P(\partial_x)P(x) = 0 \implies P(x) = 0$$

Note that this implies that if $\deg \Delta(x) = n_0$ then

$$F_V(t) = t^{n_0} F_V(1/t)$$

The same is true if V is a cone with summit a bihomogeneous polynomial $\Delta(x, y)$ of bi-degree n_x, n_y . In this case

$$F_V(q, t) = t^{n_x} q^{n_y} F_V(1/q, 1/t)$$

This explains the symmetry

$$F_{32}(q, t) = \begin{bmatrix} 5 & 11 & 9 & 4 & 1 \\ 4 & 15 & 22 & 15 & 4 \\ 1 & 4 & 9 & 11 & 5 \end{bmatrix}$$

ORBITS

Let G be a group of permutations of $1, 2, \dots, n$. Given an element $\sigma = (\sigma_1, \sigma_2, \dots, \sigma_n) \in G$ and a vector $a = (a_1, a_2, \dots, a_n)$ we set

$$\sigma a = (a_{\sigma_1}, a_{\sigma_2}, \dots, a_{\sigma_n})$$

This given, we let

$$[a]_G = \{b = \sigma a : \sigma \in G\}$$

and call it the “orbit” of a under G . The subgroup

$$G_a = \{\sigma \in G : \sigma a = a\}$$

is called the “stabilizer” of a .

If G_a is trivial then a is called “regular” and the orbit $[a]_G$ has cardinality $|G|$.

In general we have

$$|[a]_G| = |G|/|G_a|$$

In any case we set

$$J_{[a]_G} = (P : P(b) = 0 \quad \forall \quad b \in [a]_G) \quad gr J_{[a]_G} = (h(P) : P \in J_{[a]_G})$$

and

$$R_{[a]_G} = \mathbb{Q}[x_1, x_2, \dots, x_n]/J_{[a]_G}, \quad gr R_{[a]_G} = \mathbb{Q}[x_1, x_2, \dots, x_n]/gr J_{[a]_G}$$

We also set

$$H_{[a]_G} = (gr J_{[a]_G})^\perp$$

and call the elements of this space the “orbit harmonics”.

Now it follows that

$$\dim R_{[a]_G} = \dim gr R_{[a]_G} = \dim H_{[a]_G} = |G|/|G_a|$$

[next](#)

Conical Orbit Harmonics

(This is why symmetry proves the $n!$ result)

Theorem

Let a be a regular point. Set

$$n_o = \text{maxdegree } R_{[a]_G}$$

Suppose that \mathcal{B} is a basis for $R_{[a]_G}$ with

$$\mathcal{B}^{\leq n_o} = \mathcal{B} \quad \text{and} \quad |\mathcal{B}^{\leq i}| = d_o + d_1 + \cdots + d_i \quad (i = 0, 1, \dots, n_o)$$

Then the equalities

$$d_i = d_{n_o - i} \quad (i = 0, 1, \dots, n_o)$$

imply

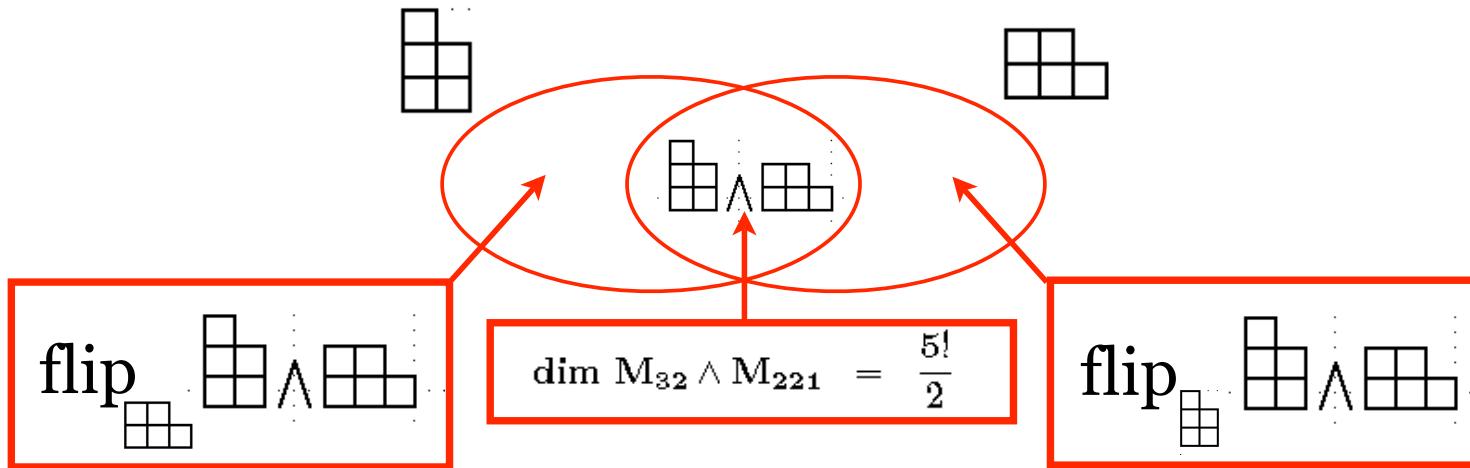
(i) $h(\mathcal{B}) = \{ h(b) : b \in \mathcal{B} \}$ is a basis for $\text{gr } R_{[a]}$

(ii) $F_{\text{gr } R_{[a]}_G}(q) = F_{H_{[a]}_G}(q) = \sum_{i=0}^{n_o} d_i q^i$,

(iii) $H_{[a]}_G$ is a cone.

(iv) The elements of $h(\mathcal{B})$ yield the desired $n!$ derivatives

Now the miracles



$$M_{221} = M_{32} \wedge M_{221} \oplus \text{flip}_{221} M_{32} \wedge M_{221}$$

$$M_{32} = M_{221} \wedge M_{32} \oplus \text{flip}_{32} M_{221} \wedge M_{32}$$

$\frac{n!}{2}$ Conjecture (still open!)

If $\alpha, \beta \vdash n$ differ only by the position of one corner cell then $\dim M_\alpha \wedge M_\beta = \frac{n!}{2}$

SF Conjecture (still open!)

If $\alpha, \beta \vdash n$ differ only by the position of one corner cell then

$$M_\alpha = M_\alpha \wedge M_\beta \oplus \text{flip}_\alpha M_\alpha \wedge M_\beta , \quad M_\beta = M_\alpha \wedge M_\beta \oplus \text{flip}_\beta M_\alpha \wedge M_\beta$$

This implies

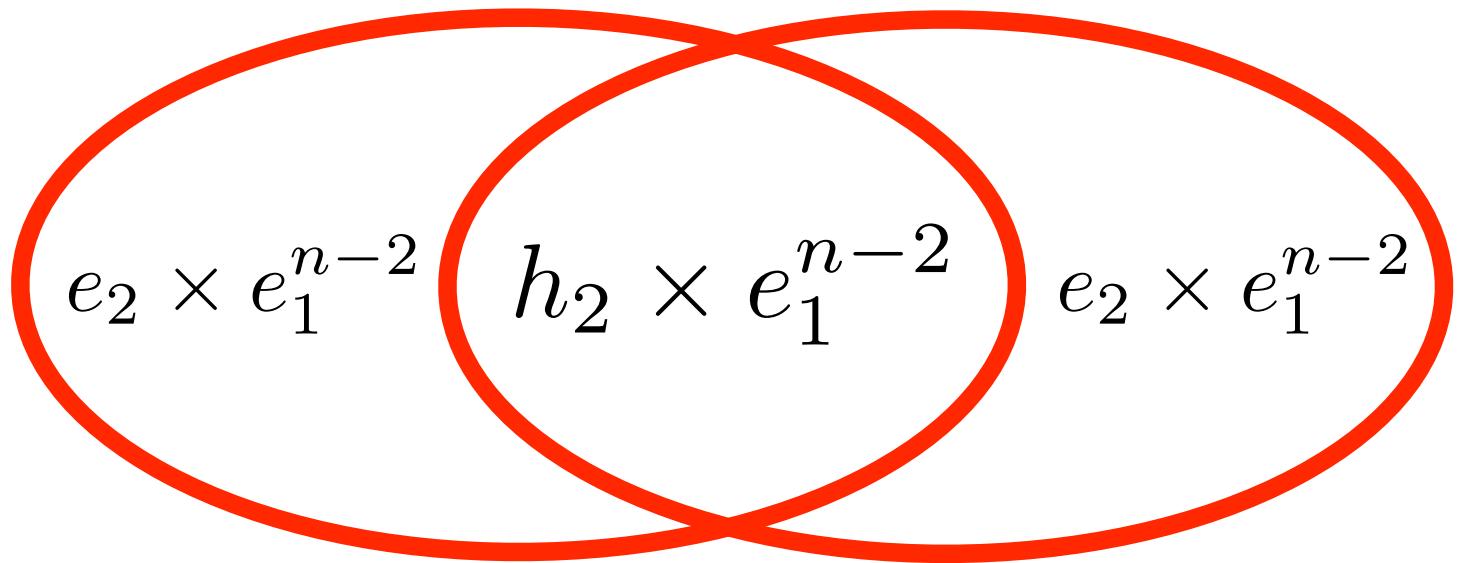
$$F_{M_\alpha \wedge M_\beta}(q, t) = \frac{T_\beta F_{M_\alpha}(q, t) - T_\alpha F_{M_\beta}(q, t)}{T_\alpha - T_\beta}$$

with $T_\mu = t^{n(\mu)} q^{n(\mu')}$ and $n(\mu) = \sum_i (i-1)\mu_i$

next

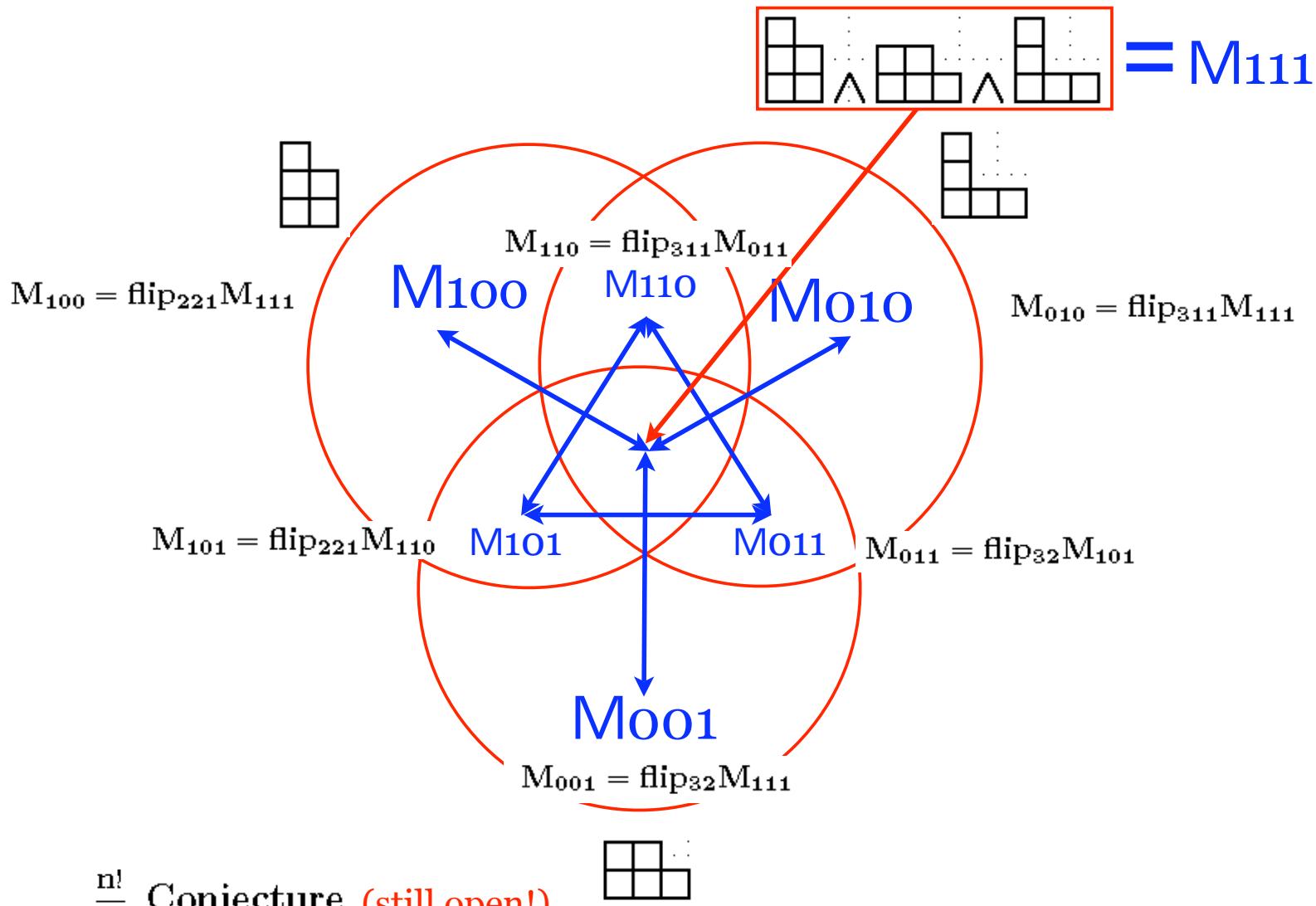
How would you split in one half
a
Left regular representation ?

$$\begin{aligned} e_1^n &= e_1^2 \times e_1^{n-2} \\ &= (h_2 + e_2) \times e_1^{n-2} \end{aligned}$$



next

More miracles



Etc Etc Etc

next

ENTER LASCOUX-LAPOINTE-MORSE k-SCHUR

The 3-Schur for 3 bounded partitions of 6

$$[1, 1, 1, 1, 1, 1], \rightarrow , t^3 s_{1, 1, 1, 1, 1, 1} + t^2 s_{2, 1, 1, 1, 1, 1} + t s_{2, 2, 1, 1, 1} + s_{2, 2, 2}$$

$$[2, 1, 1, 1, 1], \rightarrow , s_{3, 2, 1} + (s_{3, 1, 1, 1} + s_{2, 2, 1, 1}) t + t^2 s_{2, 1, 1, 1, 1}$$

$$[2, 2, 1, 1], \rightarrow , t^2 s_{2, 2, 1, 1} + t s_{3, 2, 1} + s_{3, 3}$$

$$[3, 1, 1, 1], \rightarrow , s_{4, 1, 1} + t s_{3, 1, 1, 1}$$

$$[2, 2, 2], \rightarrow , t^2 s_{2, 2, 2} + t s_{3, 2, 1} + s_{4, 2}$$

$$[3, 2, 1], \rightarrow , s_{5, 1} + (s_{4, 2} + s_{4, 1, 1}) t + t^2 s_{3, 2, 1}$$

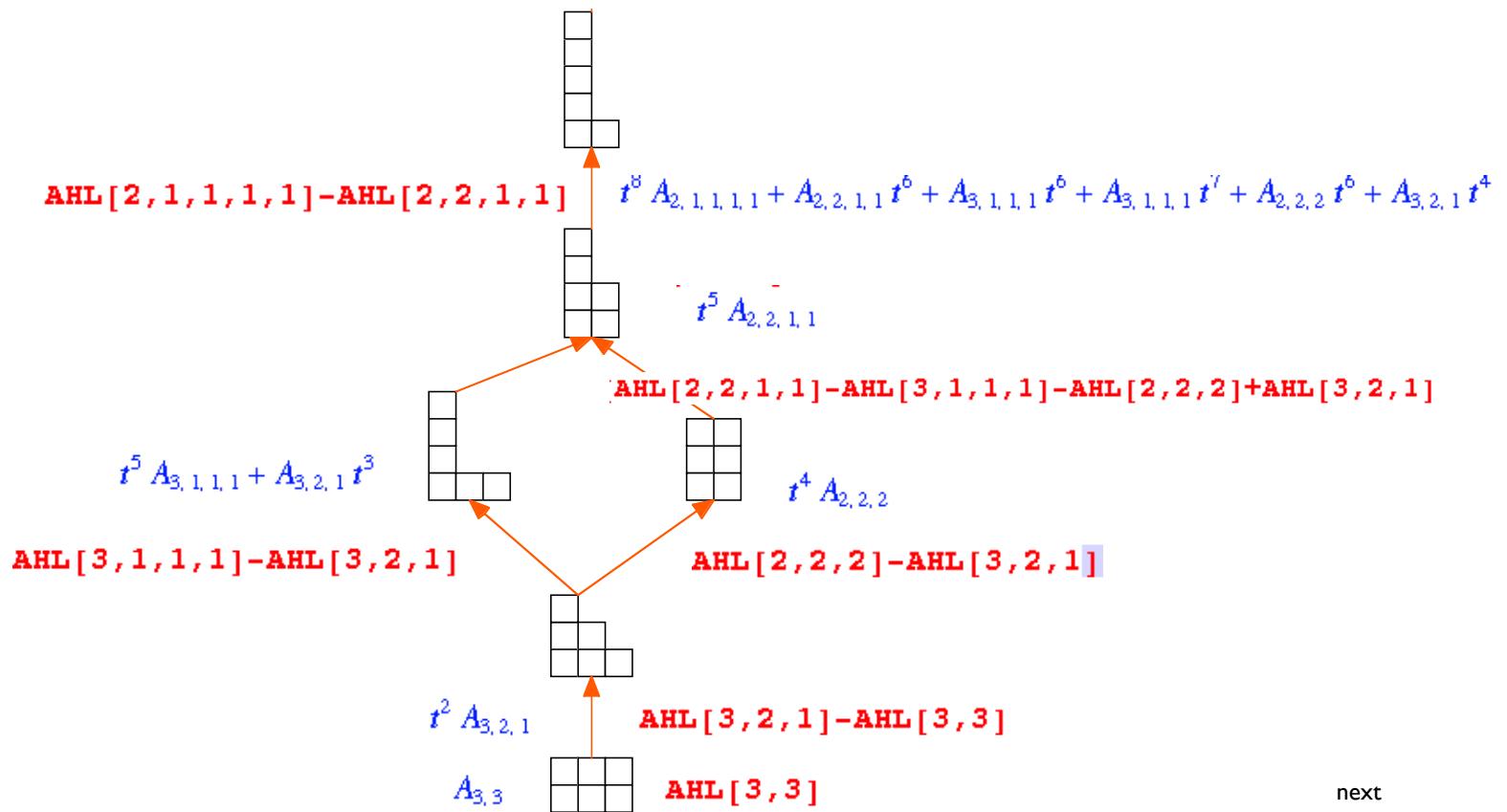
$$[3, 3], \rightarrow , s_6 + s_{5, 1} t + s_{4, 2} t^2 + s_{3, 3} t^3$$

[next](#)

Atomic Decomposition of 3-Littlewoods

(for 3-bounded partitions of 6)

(Lascoux-LaPointe-Morse 3-Schur)

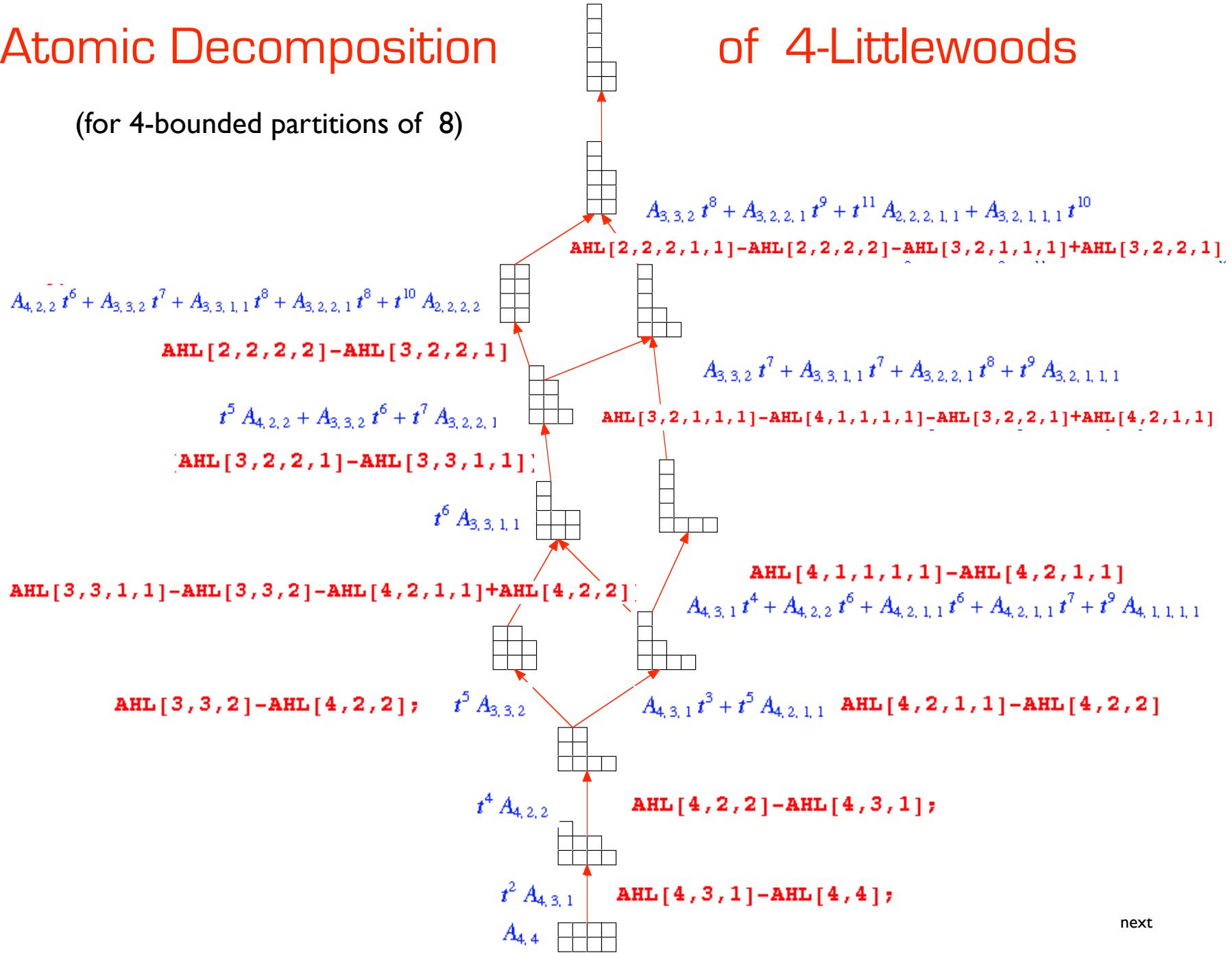


next

Atomic Decomposition

of 4-Littlewoods

(for 4-bounded partitions of 8)



next

k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4)

$$\begin{aligned}
 \text{H}([2,2]) &= q^2 A_{1,1,1,1} + (q+t) A_{2,1,1} + A_{2,2} & \begin{bmatrix} 0 & A_{2,1,1} & 0 \\ A_{2,2} & A_{2,1,1} & A_{1,1,1,1} \end{bmatrix} \\
 \text{H}([2,1,1]) &= t q A_{1,1,1,1} + (q+t^2) A_{2,1,1} + A_{2,2} & \begin{bmatrix} A_{2,1,1} & 0 \\ 0 & A_{1,1,1,1} \\ A_{2,2} & A_{2,1,1} \end{bmatrix} \\
 \text{H}([1,1,1,1]) &= t^4 A_{1,1,1,1} + (t^2+t^3) A_{2,1,1} + A_{2,2} & \begin{bmatrix} A_{1,1,1,1} \\ A_{2,1,1} \\ A_{2,1,1} \\ 0 \\ A_{2,2} \end{bmatrix}
 \end{aligned}$$

(for 2 bounded partitions of 6)

$$\begin{aligned}
 H_{2,2,2}, \quad \rightarrow, \quad , \begin{bmatrix} 0 & A_{2,2,1,1} & A_{2,1,1,1,1} & 0 \\ 0 & A_{2,2,1,1} & A_{2,1,1,1,1} & 0 \\ A_{2,2,2} & A_{2,2,1,1} & A_{2,1,1,1,1} & A_{1,1,1,1,1,1} \end{bmatrix} & H_{2,2,1,1}, \quad \rightarrow, \quad , \begin{bmatrix} A_{2,2,1,1} & A_{2,1,1,1,1} & 0 \\ 0 & A_{2,1,1,1,1} & 0 \\ 0 & A_{2,2,1,1} & A_{1,1,1,1,1,1} \\ A_{2,2,2} & A_{2,2,1,1} & A_{2,1,1,1,1} \end{bmatrix} \\
 H_{2,1,1,1,1}, \quad \rightarrow, \quad , \begin{bmatrix} A_{2,1,1,1,1} & 0 \\ 0 & 0 \\ A_{2,2,1,1} & A_{1,1,1,1,1,1} \\ A_{2,2,1,1} & A_{2,1,1,1,1} \\ 0 & A_{2,1,1,1,1} \\ 0 & 0 \\ A_{2,2,2} & A_{2,2,1,1} \end{bmatrix} & H_{1,1,1,1,1,1}, \quad \rightarrow, \quad , \begin{bmatrix} A_{1,1,1,1,1,1} \\ A_{2,1,1,1,1,1} \\ A_{2,1,1,1,1,1} \\ A_{2,1,1,1,1,1} \\ A_{2,2,1,1,1} \\ A_{2,2,1,1,1} \\ A_{2,2,1,1,1} \\ 0 \\ 0 \\ A_{2,2,2} \end{bmatrix}
 \end{aligned}$$

next

k-Schur function compatibility with Science Fiction

(for pairs of 3 bounded partitions of 5)

$$[2, 2, 1], [3, 1, 1], \rightarrow, \begin{bmatrix} A_{3, 1, 1} & A_{2, 1, 1, 1} & 0 \\ 0 & A_{2, 2, 1} & 0 \\ A_{3, 2} & A_{3, 1, 1} & A_{2, 2, 1} \end{bmatrix}$$

$$[3, 2], [3, 1, 1], \rightarrow, \begin{bmatrix} 0 & A_{2, 2, 1} & 0 & 0 \\ A_{3, 2} & A_{3, 1, 1} & A_{2, 2, 1} + A_{3, 1, 1} & A_{2, 1, 1, 1} \end{bmatrix}$$

(for pairs of 3 bounded partitions of 6)

$$[3, 2, 1], [3, 1, 1, 1], \rightarrow, \begin{bmatrix} A_{3, 2, 1} & A_{2, 2, 1, 1} + A_{3, 1, 1, 1} + A_{2, 2, 2} & A_{3, 1, 1, 1} & 0 \\ 0 & 0 & A_{2, 2, 1, 1} & A_{2, 1, 1, 1, 1} \\ A_{3, 3} & A_{3, 2, 1} & A_{2, 2, 2} + A_{3, 2, 1} & A_{3, 1, 1, 1} \end{bmatrix}$$

$$[2, 2, 2], [2, 2, 1, 1], \rightarrow, \begin{bmatrix} A_{2, 2, 2} & 0 & 0 \\ 0 & A_{2, 2, 1, 1} + A_{3, 1, 1, 1} & 0 \\ A_{3, 2, 1} & A_{2, 2, 1, 1} + A_{3, 1, 1, 1} & A_{2, 1, 1, 1, 1} \\ 0 & A_{3, 2, 1} & A_{3, 1, 1, 1} \\ A_{3, 3} & A_{3, 2, 1} & A_{2, 2, 2} \end{bmatrix}$$

$$[2, 1, 1, 1, 1], [2, 2, 1, 1], \rightarrow, \begin{bmatrix} A_{2, 2, 1, 1} + A_{3, 1, 1, 1} & 0 \\ A_{2, 2, 2} & A_{2, 1, 1, 1, 1} \\ A_{3, 2, 1} & A_{3, 1, 1, 1} \\ A_{3, 2, 1} & A_{2, 2, 1, 1} + A_{3, 1, 1, 1} + A_{2, 2, 2} \\ 0 & 0 \\ A_{3, 3} & A_{3, 2, 1} \end{bmatrix}$$

next

A new Problem

Recall that from the Haglund-Haiman-Loehr result we now have statistics $a(T), b(T)$ giving

$$F_\mu(q, t) = \sum_{T \in \text{INJ}(\mu)} t^{a(T)} q^{b(T)}$$

Jim Haglund asked for a new formula of the form

$$F_\mu(q, t) = \sum_{T \in \text{ST}(\mu)} P_T(q, t) \quad (P_T(q, t) \in \mathbb{N}[q, t])$$

In a new burst of genius Jim conjectured that

$$F_{2^b 1^{a-b}}(q, t) = \sum_{T \in \text{ST}(2^b 1^{a-b})} \prod_{i \in T} (1 + t^{\text{drop}_T(i)}) \prod_{i \in \text{col}_2(T)} (t^{\text{NW} > i} + q) \quad (*)$$

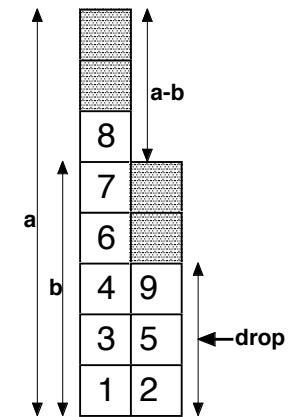
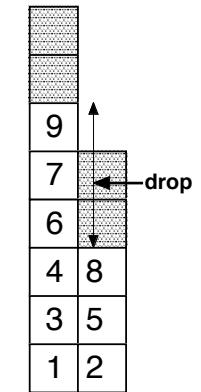
With “ $\text{drop}_T(i)$ ” is the number of rows in $T_{\leq i}$ of the same length as the row that contains i , and “ $\text{NW} > i$ ” gives the number of entries in T strictly NW of i that are larger than i .

This is equivalent to the recursion

$$F_{2^b 1^{a-b}}(q, t) = (1 + q)[b]_t F_{2^{b-1} 1^{a-b+1}}(q, t) + [a - b]_t t^b F_{2^b 1^{a-b-1}}(q/t, t) \quad (**)$$

This is quite remarkable since in the 20 years since the Macdonald conjectures, we never suspected that we could have a purely combinatorial recursion for $F_\mu(q, t)$.

It is a challenging problem to find the general form of $(*)$ and $(**)$



next

Proof of the Haglund conjecture

Science Fiction suggested that the recursion

$$F_{2^{b_1}a-b}(q, t) = (1+q)[b]_t F_{2^{b-1}1^{a-b+1}}(q, t) + [a-b]_t t^b F_{2^{b_1}a-b-1}(q/t, t)$$

is none other than the result of applying $\partial_{p_1}^{n-1}$ to the Macdonald polynomial identity

$$\partial_{p_1} \tilde{H}_{2^{b_1}a-b}(X; q, t) = (1+q)[b]_t \tilde{H}_{2^{b-1}1^{a-b+1}}(X; q, t) + [a-b]_t t^b \tilde{H}_{2^{b_1}a-b-1}(X; q/t, t)$$

for instance when $a = 3$ and $b = 2$ this is

$$\partial_{p_1} \begin{smallmatrix} & 1 \\ & 1 \\ 2 & 1 \end{smallmatrix}(q, t) = (1+t)(1+q) \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}(q, t) + t^2 \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}(q/t, t)$$

Now Science Fiction suggests writing

$$\begin{smallmatrix} & 1 \\ & 1 \\ 2 & 1 \end{smallmatrix} = \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} + t^3 q \downarrow \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}$$

$$\begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} = \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} + t^2 q^2 \downarrow \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}$$

thus

$$\begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} = \frac{q \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \\ 1 & 1 \end{smallmatrix} - t \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}}{q-t} =_{\text{def}} \phi_{32}$$

$$t^3 q \downarrow \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} = t \frac{\begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \\ 1 & 1 \end{smallmatrix} + \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}}{q-t} =_{\text{def}} t \psi_{32}$$

$$t^2 q^2 \downarrow \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} \wedge \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} = q \frac{\begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \\ 1 & 1 \end{smallmatrix} - \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \\ 1 & 1 \end{smallmatrix}}{q-t} =_{\text{def}} q \psi_{32}$$

and thus

$$\begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} = \phi_{32} + t \psi_{32}$$

$$\begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix} = \phi_{32} + q \psi_{32}$$

Now in terms of 2-Schur we get

$$\phi_{32} = q A_{211} + A_{22}, \quad q \psi_{32} = q^2 A_{1111} + qt A_{211}$$

this gives

$$t^2 \begin{smallmatrix} & 1 \\ & 1 \\ 1 & 1 \end{smallmatrix}(q/t, t) = \begin{bmatrix} A_{22} & A_{211} & 0 \\ 0 & A_{211} & 0 \\ 0 & 0 & A_{1111} \end{bmatrix} = t^2 \phi_{32} + q \psi_{32}$$

and using this in (*) we are reduced to show that

$$\partial_{p_1} \begin{smallmatrix} & 1 \\ & 1 \\ 2 & 1 \end{smallmatrix}(q, t) = (1+t)(1+q)(\phi_{32} + t \psi_{32}) + t^2 \phi_{32} + q \psi_{32}$$

next

Using “pistol” Macdonalds

$$\partial_{p_1} = t^b [a-b]_t + q[b]_t + [b]_t$$

$\partial_{p_1} = [a-b]_t + t^{a-b}[b]_t + q[b]_t$

$$\partial_{p_1} = \frac{(t^{a-b}-1)}{t^{a-b}-q} + (1-q)$$

$\phi_{ab} = \square \wedge \square$

$$\partial_{p_1} = t^b [a-b]_t (\phi_{ab} + q\psi_{ab}) + q[b]_t (\phi_{ab} + \psi_{ab}) + [b]_t (\phi_{ab} + t^{a-b}\psi_{ab})$$

$$(1+q)[b]_t (\phi_{ab} + t^{a-b}\psi_{ab}) + [a-b]_t (t^b\phi_{ab} + q\psi_{ab})$$

while my extension of the Haglund recursion is

$$\partial_{p_1} = (1+q)[b]_t (\phi_{ab} + t^{a-b}\psi_{ab}) + [a-b]_t t^b (q/t, t)$$

so we are reduced to proving the identity

$$t^b (q/t, t) = t^b \phi_{ab} + q\psi_{ab}$$

next

we are still Searching for “gistol” Macdonalds

(1) A “gistol” is a lattice diagram that can be transformed to a skew diagram by row and column interchanges

(2) We postulate the existence of a family of polynomials indexed by gistsols with the following basic properties:

$$\left\{ \begin{array}{ll} (0) & G_D(x; q, t) = H_\mu(x; q, t) \quad \text{if } D \text{ is the diagram of } \mu \\ (1) & G_{D_1}(x; q, t) = G_{D_2}(x; q, t) \quad \text{if } D_1 \approx D_2 \\ (2) & G_{D_1}(x; q, t) = G_{D_2}(x; t, q) \quad \text{if } D_2 \approx D'_1 \\ (3) & G_D(x; q, t) = G_{D_1}(x; q, t)G_{D_2}(x; q, t) \quad \text{if } D \approx D_1 \times D_2 \\ (4) & \partial_{p_1} G_D(x; q, t) = \sum_{s \in D} w_{s,D}(q, t) G_{D/s}(x; q, t) , \end{array} \right.$$

Representation theoretical reasons suggest that,

in the case that D is a skew diagram,

a) $w_1[s, D] = t^{l_D(s)} q^{a'_D(s)}$ and b) $w_2[s, D] = t^{l'_D(s)} q^{a_D(s)}$

Note: these properties overdetermine the family $\{G_D(x; q, t)\}_D$,

existence is by no means guaranteed.

next

THE END